

```
> restart;
with(plots) :
with(LinearAlgebra) :
with(IntegrationTools) :
```

Finite element methods in one spatial dimension

The purpose of this worksheet is to develop a finite element algorithm to solve partial differential equations of the form

$$\rho(x) \frac{\partial^2 u}{\partial t^2} + \lambda(x) \frac{\partial u}{\partial t} = \frac{\partial}{\partial x} \left[P(x) \frac{\partial u}{\partial x} \right] + Q(x)u - R(x), \quad u(t, -1) = \alpha, \quad u(t, 1) = \beta.$$

In these equations, $\{\rho, \lambda, P, Q, R\}$ are known functions of x while $\{\alpha, \beta\}$ are known constants.

Weak form of the problem

Here is the PDE we want to solve:

```
> pde := rho(x)*Diff(u(t,x),t,t) + lambda(x)*Diff(u(t,x),t) - Diff(P(x)*Diff(u(t,x),x),x) -
Q(x)*u(t,x) + R(x);
```

$$pde := \rho(x) \left(\frac{\partial^2}{\partial t^2} u(t,x) \right) + \lambda(x) \left(\frac{\partial}{\partial t} u(t,x) \right) - \left(\frac{\partial}{\partial x} \left(P(x) \left(\frac{\partial}{\partial x} u(t,x) \right) \right) \right) - Q(x) u(t,x) + R(x) \quad (1.1)$$

We do not want Maple to expand out the derivatives in the third term, so we use the inert differentiation command **Diff**. The first step in the finite element solution of this PDE is deriving the so-called "weak form" of the equation. This involves multiplying the PDE by a test function of x , which we denote by $\phi_i(x)$, and integrating over $[-1, 1]$. Note that right now $\phi_i(x)$ is an arbitrary function, but in the next section we will identify it with one member of an N -dimensional "basis": $\{\phi_i(x)\}_{i=1}^N$.

```
> eq1 := int(pde*phi[i](x),x=-1..1);
eq2 := Expand(eq1);
```

$$eq1 := \int_{-1}^1 \left(\rho(x) \left(\frac{\partial^2}{\partial t^2} u(t,x) \right) + \lambda(x) \left(\frac{\partial}{\partial t} u(t,x) \right) - \left(\frac{\partial}{\partial x} \left(P(x) \left(\frac{\partial}{\partial x} u(t,x) \right) \right) \right) - Q(x) u(t,x) + R(x) \right) \phi_i(x) dx$$

$$eq2 := \int_{-1}^1 \phi_i(x) \rho(x) \left(\frac{\partial^2}{\partial t^2} u(t,x) \right) dx + \int_{-1}^1 \phi_i(x) \lambda(x) \left(\frac{\partial}{\partial t} u(t,x) \right) dx - \left(\int_{-1}^1 \phi_i(x) \left(\frac{\partial}{\partial x} \left(P(x) \left(\frac{\partial}{\partial x} u(t,x) \right) \right) \right) \right) \quad (1.2)$$

$$dx) - \left(\int_{-1}^1 \phi_i(x) Q(x) u(t, x) dx \right) + \int_{-1}^1 \phi_i(x) R(x) dx$$

The third term in (1.2) can be integrated by parts:

> eq3 := applyop(u->Parts(u, phi[i](x)), 3, eq2);

$$\begin{aligned} eq3 := & \int_{-1}^1 \phi_i(x) \rho(x) \left(\frac{\partial^2}{\partial t^2} u(t, x) \right) dx + \int_{-1}^1 \phi_i(x) \lambda(x) \left(\frac{\partial}{\partial t} u(t, x) \right) dx - P(1) D_2(u)(t, 1) \phi_i(1) + P(-1) D_2(u)(t, \\ & -1) \phi_i(-1) + \int_{-1}^1 P(x) \left(\frac{\partial}{\partial x} u(t, x) \right) \left(\frac{d}{dx} \phi_i(x) \right) dx - \left(\int_{-1}^1 \phi_i(x) Q(x) u(t, x) dx \right) + \int_{-1}^1 \phi_i(x) R(x) dx \end{aligned} \quad (1.3)$$

Now, let us assume that our test function $\phi_i(x)$ vanishes at either end of the region of integration:

> eq4 := [phi[i](-1)=0, phi[i](1)=0];

eq5 := Expand(convert(subs(eq4, eq3), diff));

$$eq4 := [\phi_i(-1) = 0, \phi_i(1) = 0]$$

$$\begin{aligned} eq5 := & \int_{-1}^1 \phi_i(x) \rho(x) \left(\frac{\partial^2}{\partial t^2} u(t, x) \right) dx + \int_{-1}^1 \phi_i(x) \lambda(x) \left(\frac{\partial}{\partial t} u(t, x) \right) dx + \int_{-1}^1 P(x) \left(\frac{\partial}{\partial x} u(t, x) \right) \left(\frac{d}{dx} \phi_i(x) \right) dx - \left(\int_{-1}^1 \phi_i(x) Q(x) u(t, x) dx \right) + \int_{-1}^1 \phi_i(x) R(x) dx \end{aligned} \quad (1.4)$$

This is the "weak form" of the original PDE (1.1), and it is the central equation of our finite element analysis. Our goal will be to find a solutions of (1.4) for a particular choice of N basis functions $\{\phi_i(x)\}_{i=1}^N$; i.e., (1.4) represents one of a system of N equations. In what sense is equation (1.4) weak? For a given choice of basis functions, there will be a solution $u(t, x)$ that satisfies the associated system of equations, yet are not solutions of the PDE (1.1). However, for a reasonable choice of basis, the hope is that the solutions of (1.4) are good approximations to the solutions of (1.1). This is the essence of the finite element analysis, we seek exact solutions of the weak equations that approximate the solutions of the "strong" equations (1.1).

▼ Basis functions

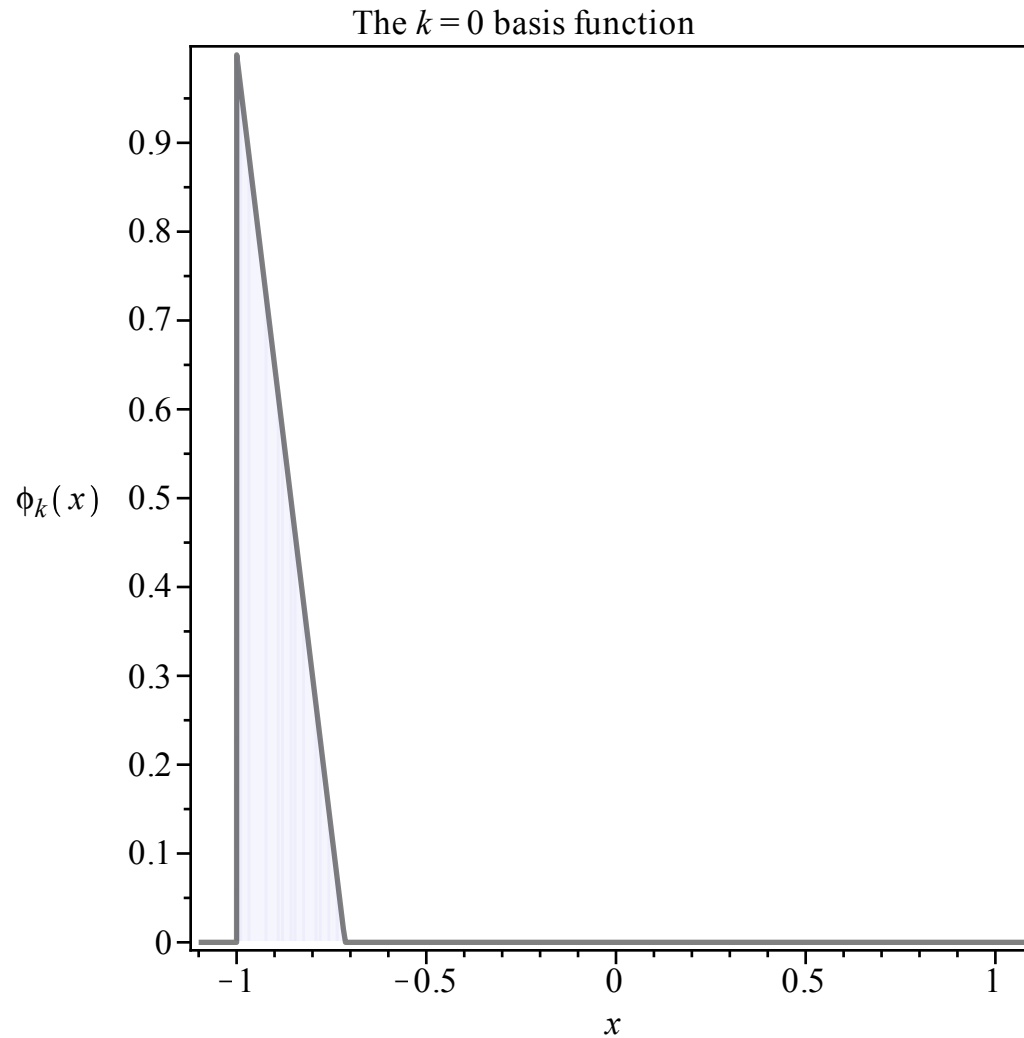
We cannot proceed any further until we choose a basis. We will restrict ourselves to simplest choice: piecewise linear functions. (Of course, many other choices of basis are possible.) To be more specific, we introduce a lattice on the interval $[-1, 1]$ of the form

$x_i = -1 + \frac{2i}{N+1}$, with $i = 0 \dots N$. With this choice, $x_0 = -1$ and $x_{N+1} = 1$. A basis function will be a piecewise linear function defined on $[-1, 1]$ whose value at each position on the lattice is

$$\phi_i(x_j) = \begin{cases} 0 & x_j = x_i \\ 1 & x_j \neq x_i \end{cases}$$

Here is some code that defines the basis for a particular choice of N and plots the functions in a movie. Notice the use of **piecewise** to define each member of the basis.

```
> phi_def := 'phi_def':
N := 6:
X := i -> -1 + 2*i/(N+1):
phi_def[0] := phi[0](x) = piecewise(x>X(0) and x < X(1), (X(1)-x)/(X(1)-X(0))):
for i from 1 to N do:
  phi_def[i] := phi[i](x) = piecewise(x>X(i-1) and x <= X(i), (x-X(i-1))/(X(i)-X(i-1)),
x>X(i) and x < X(i+1), (X(i+1)-x)/(X(i+1)-X(i))):
od:
phi_def[N+1] := phi[N+1](x) = piecewise(x>X(N) and x < X(N+1), (x-X(N))/(X(N+1)-X(N))):
phi_def := convert(phi_def,list):
i := 'i':
display([seq(plot([rhs(phi_def[i])$2],x=-1.1..1.1,color=[black,"Lavender"],thickness=[2,
0],transparency=0.5,filled=[false,true],axes=boxed,labels=[x,typeset(phi[k](x))],title=
typeset("The ",k=i-1," basis function")),i=1..N+2)],insequence=true);
```



These basis functions are the "finite elements" we will use to construct the numeric solution to **(1.1)**. Note that the basis can be used to obtain a piecewise linear approximation to a function $f(x)$ defined on $[-1, 1]$:

$$f(x) \approx \sum_{i=0}^N C_i \phi_i(x), \quad C_i = f(x_i).$$

The LHS will equal the RHS when at the lattice points (i.e., when $x = x_i$). In between lattice points, the RHS will be a continuous linear function. Here are some plots showing how the decomposition works:

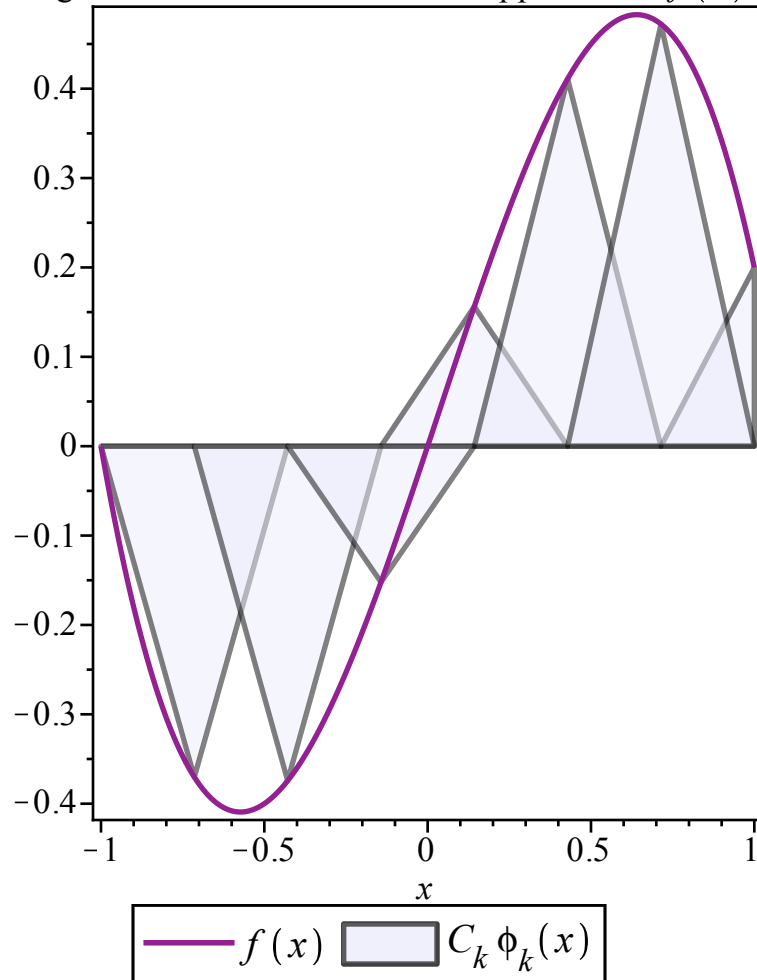
```

> f := x -> x*(11/10-x)*(1+x);
c := Array(0..N+1,[seq(f(X(i)),i=0..N+2)]):
Data1 := [[X(0),0],[X(1),0],[X(0),c[0]],seq([X(i-1),0],[X(i+1),0],[X(i),c[i]],i=1..N),
[X(N),0],[X(N+1),0],[X(N+1),c[N+1]]]:
Data2 := subs(phi_def,add(c[i]*phi[i](x),i=0..N+1)):
Data3 := [[-1,0],seq([X(i),c[i]],i=0..N+1),[1,0]]:
p1 := plot(f(x),x=-1..1,legend=['''f(x)'''],thickness=2,color="Purple"):
p2 := polygonplot(Data1,color="Lavender",axes=boxed,transparency=0.5,thickness=2,legend=
typeset(C[k]*phi[k](x))):
p3 := plot([Data2],x=-1..1,color=black,axes=boxed,legend=[typeset(Sum(C[k]*phi[k](x),k=0..
N+1))],transparency=0.5,thickness=2):
p4 := polygonplot(Data3,color="Lavender",transparency=0.5):
display(Array([display([p1,p2],title=typeset("Weighted basis functions used to approximate
",'f(x)'))),display([p1,p3,p4],title=typeset("Piecewise linear approximation to",'f(x)
')))]));

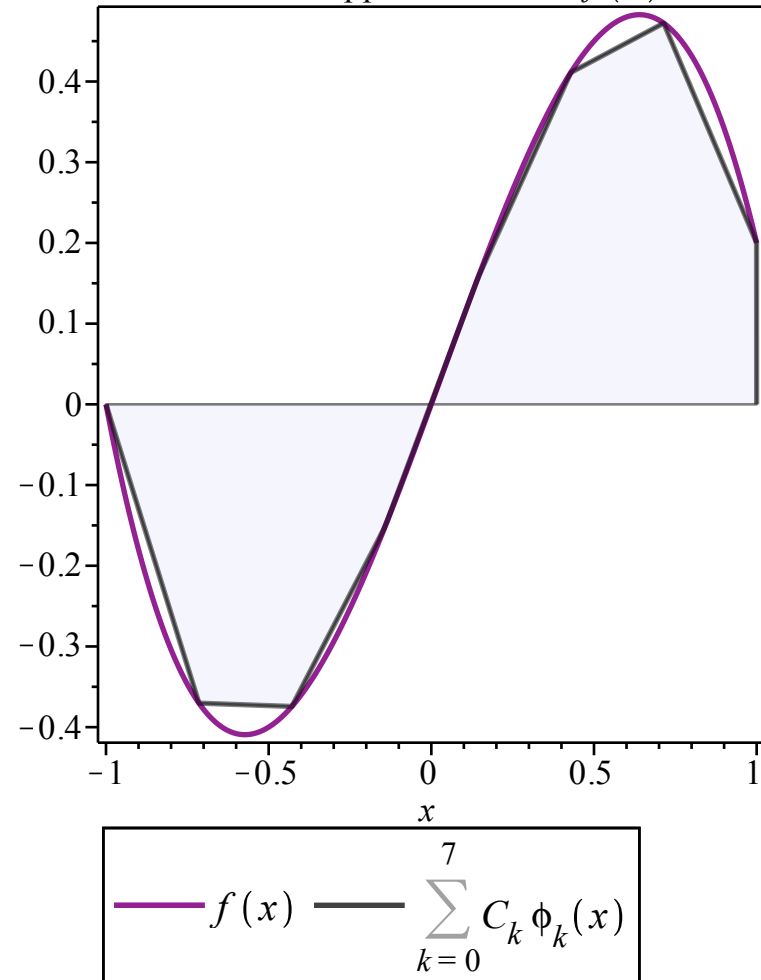
```

$$f := x \rightarrow x \left(\frac{11}{10} - x \right) (x + 1)$$

Weighted basis functions used to approximate $f(x)$



Piecewise linear approximation to $f(x)$



Now, let's us write down a piewise linear approximation to the solution of the PDE (1.1) using these basis functions:

```
> N := 'N':
eq6 := u(t,x) = alpha*phi[0](x) + Sum(a[j](t)*phi[j](x), j=1..N) + beta*phi[N+1](x);
```

$$eq6 := u(t, x) = \alpha \phi_0(x) + \sum_{j=1}^N a_j(t) \phi_j(x) + \beta \phi_{N+1}(x)$$

(2.1)

Notice that we have enforced the boundary conditions $u(t, -1) = \alpha$ and $u(t, +1) = \beta$ by forcing the coefficients of $\phi_0(x)$ and $\phi_{N+1}(x)$ to be α and β , respectively. Also note that all the time dependence of the solution is carried by the other coefficients $a_j(t)$. Our goal will be to solve for these coefficients.

Matrix form

We now substitute the decomposition of $u(t, x)$ in terms of basis functions (2.1) into the weak form of the PDE (1.4):

`> eq7 := Expand(convert(subs(eq6, eq5), diff));`

$$\begin{aligned}
 \text{eq7} := & \int_{-1}^1 \phi_i(x) \rho(x) \left(\sum_{j=1}^N \left(\frac{d^2}{dt^2} a_j(t) \right) \phi_j(x) \right) dx + \int_{-1}^1 \phi_i(x) \lambda(x) \left(\sum_{j=1}^N \left(\frac{d}{dt} a_j(t) \right) \phi_j(x) \right) dx + \alpha \left(\right. \\
 & \left. \int_{-1}^1 P(x) \left(\frac{d}{dx} \phi_i(x) \right) \left(\frac{d}{dx} \phi_0(x) \right) dx \right) + \int_{-1}^1 P(x) \left(\frac{d}{dx} \phi_i(x) \right) \left(\sum_{j=1}^N a_j(t) \left(\frac{d}{dx} \phi_j(x) \right) \right) dx + \beta \left(\right. \\
 & \left. \int_{-1}^1 P(x) \left(\frac{d}{dx} \phi_i(x) \right) \left(\frac{d}{dx} \phi_{N+1}(x) \right) dx \right) - \alpha \left(\int_{-1}^1 \phi_i(x) Q(x) \phi_0(x) dx \right) - \left(\int_{-1}^1 \phi_i(x) Q(x) \left(\sum_{j=1}^N a_j(t) \phi_j(x) \right) \right. \\
 & \left. dx \right) - \beta \left(\int_{-1}^1 \phi_i(x) Q(x) \phi_{N+1}(x) dx \right) + \int_{-1}^1 \phi_i(x) R(x) dx
 \end{aligned} \tag{3.1}$$

We now evaluate this expression for each of the $\phi_i(x)$ basis functions with an unknown coefficient; i.e., for $i = 1 \dots N$. It is not difficult to see that this yields a *matrix* ODE for the unknown coefficients:

$$M \frac{d^2 \mathbf{a}}{dt^2} + C \frac{d\mathbf{a}}{dt} + K\mathbf{a} = \mathbf{f}, \quad \mathbf{a} = \begin{bmatrix} a_1(t) \\ \vdots \\ a_N(t) \end{bmatrix}$$

$$M = \begin{bmatrix} m_{1,1} & \cdots & m_{1,N} \\ \vdots & & \vdots \\ m_{N,1} & \cdots & m_{N,N} \end{bmatrix}, \quad C = \begin{bmatrix} c_{1,1} & \cdots & c_{1,N} \\ \vdots & & \vdots \\ c_{N,1} & \cdots & c_{N,N} \end{bmatrix}, \quad K = \begin{bmatrix} k_{1,1} & \cdots & k_{1,N} \\ \vdots & & \vdots \\ k_{N,1} & \cdots & k_{N,N} \end{bmatrix}, \quad \mathbf{f} = \begin{bmatrix} F_1 \\ \vdots \\ F_N \end{bmatrix},$$

where:

$$m_{i,j} = \int_{-1}^1 \phi_i(x) \rho(x) \phi_j(x) dx, \quad c_{i,j} = \int_{-1}^1 \phi_i(x) \lambda(x) \phi_j(x) dx, \quad k_{i,j} = \int_{-1}^1 \left\{ \left[\frac{d}{dx} \phi_i(x) \right] P(x) \left[\frac{d}{dx} \phi_j(x) \right] - \phi_i(x) Q(x) \phi_j(x) \right\} dx,$$

$$F_i = -\alpha k_{i,0} - \beta k_{i,N+1} - \int_{-1}^1 \phi_i(x) R(x) dx$$

Finite element methods were first used extensively in structural engineering, so the names of the above objects are inherited from that field: M is the mass matrix, C is the damping matrix, K is the stiffness matrix, and \mathbf{f} is the applied force. The three matrices are symmetric, and since the basis functions $\phi_i(x)$ are only non-zero in the interval $[x_{i-1}, x_{i+1}]$ they will also be tridiagonal. Here are some procedures we can use to generate the elements of these objects:

```
> # this procedure generates an array of the basis function written as mappings (its output is used in the following):
GenerateBasis := proc(N)
    local X, phi, i:
    X := i -> -1 + 2*i/(N+1):
    phi := Array(0..N+1):
    phi[0] := unapply(piecewise(x>X(0) and x < X(1), (X(1)-x)/(X(1)-X(0))), x):
    for i from 1 to N do:
        phi[i] := unapply(piecewise(x>X(i-1) and x <= X(i), (x-X(i-1))/(X(i)-X(i-1)), x>X
(i) and x < X(i+1), (X(i+1)-x)/(X(i+1)-X(i))), x):
    od:
    phi[N+1] := unapply(piecewise(x>X(N) and x < X(N+1), (x-X(N))/(X(N+1)-X(N))), x):
    phi:
end proc:

# this procedure generates elements of the mass matrix given the basis functions phi
m := proc(phi, rho, i, j):
    int(phi[i](x)*rho(x)*phi[j](x), x=-1..1):
end proc:
```



```

# this procedure generates elements of the damping matrix given the basis functions phi
c := proc (phi, lambda, i, j) :
    int(phi[i](x)*lambda(x)*phi[j](x), x=-1..1) :
end proc:

# this procedure generates elements of the stiffness matrix given the basis functions phi
k := proc (phi, P, Q, i, j) :
    int(D(phi[i])(x)*P(x)*D(phi[j])(x)-phi[i](x)*Q(x)*phi[j](x), x=-1..1) :
end proc:

# this procedure generates components of the applied force given the basis functions phi
F := proc (phi, P, Q, R, i, alpha, beta) :
    -alpha*k(phi, P, Q, i, 0)-beta*k(phi, P, Q, i, rhs(ArrayDims(phi))) -int(phi[i](x)*R(x), x=-1.
.1) :
end proc:

```

Here is an example of how the matrices and applied force look for a particular choice of parameters:

```

> Phi := GenerateBasis(4) :
N := rhs(ArrayDims(Phi))-1;
l := 1;
rho := x -> 1;
lambda := x -> x;
P := x -> (1-x^2);
Q := x -> 1*(1+1);
R := x -> 4;
alpha := (-1)^1;
beta := 1;
M := Matrix([seq([seq(m(Phi, rho, i, j), i=1..N)], j=1..N)]):
C := Matrix([seq([seq(c(Phi, lambda, i, j), i=1..N)], j=1..N)]):
K := Matrix([seq([seq(k(Phi, P, Q, i, j), i=1..N)], j=1..N)]):
f := Vector([seq(F(Phi, P, Q, R, i, alpha, beta, N), i=1..N)]):
pde;
'M'=M, 'C'=C, 'K'=K, 'f'=f;

```

$N := 4$

$l := 1$

$\rho := x \rightarrow 1$

$\lambda := x \rightarrow x$

$$\begin{aligned}
 P &:= x \rightarrow 1 - x^2 \\
 Q &:= x \rightarrow l(l+1) \\
 R &:= x \rightarrow 4 \\
 \alpha &:= -1 \\
 \beta &:= 1
 \end{aligned}$$

$$\frac{\partial^2}{\partial t^2} u(t, x) + x \left(\frac{\partial}{\partial t} u(t, x) \right) - \left(\frac{\partial}{\partial x} \left((1 - x^2) \left(\frac{\partial}{\partial x} u(t, x) \right) \right) \right) - 2u(t, x) + 4$$

$$M = \begin{bmatrix} \frac{4}{15} & \frac{1}{15} & 0 & 0 \\ \frac{1}{15} & \frac{4}{15} & \frac{1}{15} & 0 \\ 0 & \frac{1}{15} & \frac{4}{15} & \frac{1}{15} \\ 0 & 0 & \frac{1}{15} & \frac{4}{15} \end{bmatrix}, C = \begin{bmatrix} -\frac{4}{25} & -\frac{2}{75} & 0 & 0 \\ -\frac{2}{75} & -\frac{4}{75} & 0 & 0 \\ 0 & 0 & \frac{4}{75} & \frac{2}{75} \\ 0 & 0 & \frac{2}{75} & \frac{4}{25} \end{bmatrix}, K = \begin{bmatrix} \frac{12}{5} & -\frac{11}{5} & 0 & 0 \\ -\frac{11}{5} & 4 & -\frac{13}{5} & 0 \\ 0 & -\frac{13}{5} & 4 & -\frac{11}{5} \\ 0 & 0 & -\frac{11}{5} & \frac{12}{5} \end{bmatrix}, f = \begin{bmatrix} -\frac{13}{5} \\ -\frac{8}{5} \\ -\frac{8}{5} \\ -\frac{3}{5} \end{bmatrix}$$

(3.2)

As claimed above, the matrices are symmetric and tridiagonal, which implies that the above method is not the most efficient way of calculating them. The following procedures are faster:

```

> mass := proc(phi, rho)
  local N, M, i:
  N := rhs(ArrayDims(phi))-1:
  M := Matrix(N,N,shape=symmetric,datatype=float):
  for i from 1 to N do:
    M[i,i] := evalf(m(phi, rho, i, i)):
    if (i<>N) then M[i,i+1] := evalf(m(phi, rho, i, i+1)) fi:
  od:
  M:
end proc:

damping := proc(phi, lambda)
  local N, C, i:
  N := rhs(ArrayDims(phi))-1:
  C := Matrix(N,N,shape=symmetric,datatype=float):

```

```

    for i from 1 to N do:
      C[i,i] := evalf(c(phi,lambda,i,i)):
      if (i<>N) then C[i,i+1] := evalf(c(phi,lambda,i,i+1)) fi:
    od:
  C;
end proc:

stiffness := proc(phi,P,Q)
  local N, K, i:
  N := rhs(ArrayDims(phi))-1:
  K := Matrix(N,N,shape=symmetric,datatype=float):
  for i from 1 to N do:
    K[i,i] := evalf(k(phi,P,Q,i,i)):
    if (i<>N) then K[i,i+1] := evalf(k(phi,P,Q,i,i+1)) fi:
  od:
  K;
end proc:

force := proc(phi,P,Q,R,alpha,beta)
  local N, K, i:
  N := rhs(ArrayDims(phi))-1:
  Vector([seq(evalf(F(phi,P,Q,R,i,alpha,beta,N)),i=1..N)]):
end proc:

```

▼ Elliptic boundary value problems

In this section, we will concentrate on the time-independent case when $\rho(x) = \lambda(x) = 0$ in **(1.1)**; i.e., the equation we want to solve is:

$$0 = \frac{d}{dx} \left[P(x) \frac{du}{dx} \right] + Q(x)u - R(x), \quad u(-1) = \alpha, \quad u(1) = \beta.$$

(We have re-written $u(t, x) \mapsto u(x)$.) This is an ODE boundary value problem. In this case, the matrix equation we need to solve for the amplitudes \mathbf{a} of the basis functions is

$$K\mathbf{a} = \mathbf{f}.$$

The following procedure calculates the coefficients and plots the finite element solution:

```

> EllipticSolver := proc(N,P,Q,R,alpha,beta)
  local Phi,K,i,f,ans:

```

```

Phi := GenerateBasis(N) :
K := stiffness(Phi,P,Q) :
f := force(Phi,P,Q,R,alpha,beta) :
ans := LinearSolve(K,f) ;
ans := Array(0..N+1,[alpha,seq(ans[i],i=1..N),beta]) :
plot(add(ans[i]*Phi[i](x),i=0..N+1),x=-1..1,color="Lavender",filled=true,legend="Finite
element solution",axes=boxed) :
end proc:

```

We illustrate the output of the above code for the following example with exact solution **u_sol**. (This is actually Legendre's differential equation whose solutions are the Legendre functions $P_l(x)$.)

```

> l := 3:
P := x -> (1-x^2) :
Q := x -> 1*(1+1) :
R := x -> 0:
alpha := (-1)^1;
beta := 1;
rho := x -> 0:
lambda := x -> 0:
ode := subs(u(t,x)=u(x),pde) ;
u_sol := subs(dsolve([ode,u(-1)=alpha,u(1)=beta]),u(x)) ;
ans := EllipticSolver(10,P,Q,R,alpha,beta) :
display([ans,plot(u_sol,x=-1..1,color="Purple",legend="exact solution")],axes=boxed,
labels=[x,u(x)]) ;

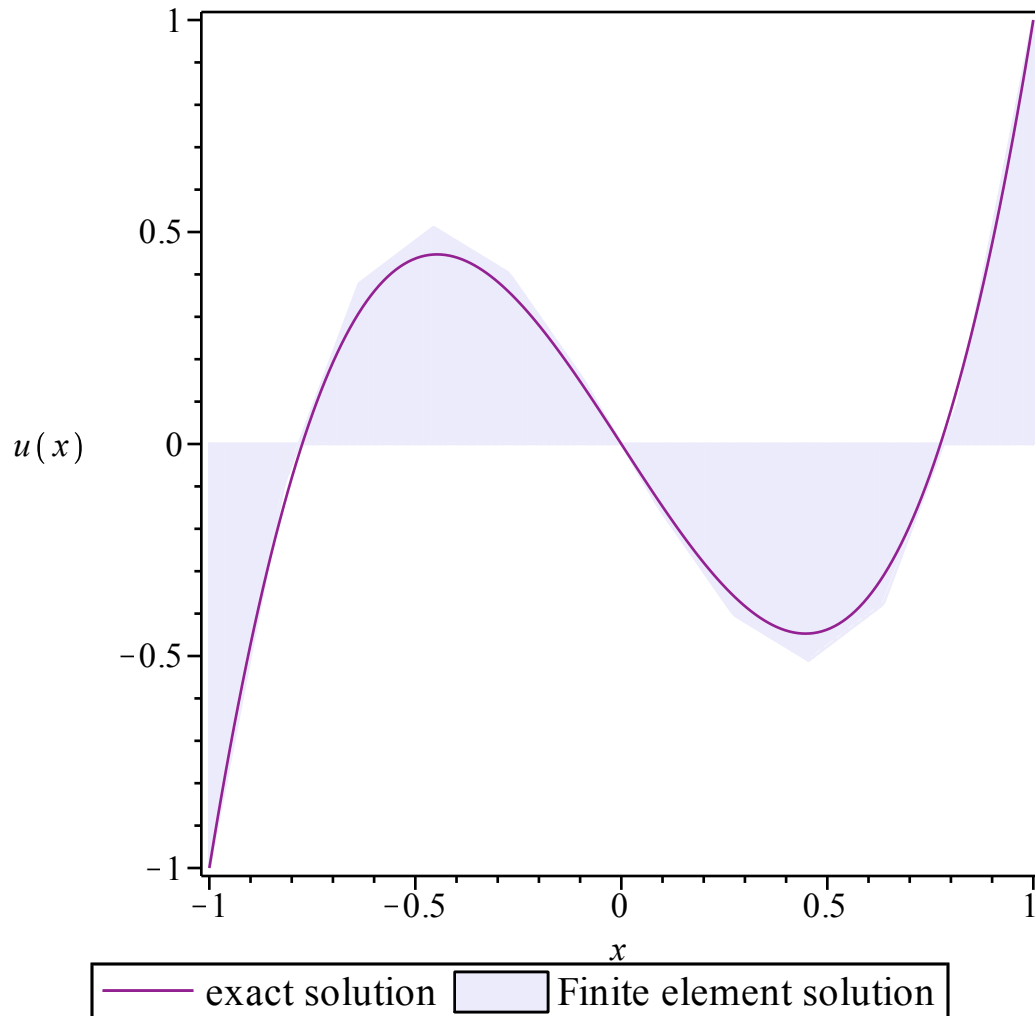
```

$$\alpha := -1$$

$$\beta := 1$$

$$ode := - \left(\frac{d}{dx} \left((1-x^2) \left(\frac{d}{dx} u(x) \right) \right) \right) - 12 u(x)$$

$$u_sol := -\frac{3}{2}x + \frac{5}{2}x^3$$



▼ Parabolic initial value problems

In this section, we will concentrate on case when $\rho(x) = 0$ in **(1.1)**; i.e., the equation we want to solve is:

$$\lambda(x) \frac{\partial u}{\partial t} = \frac{\partial}{\partial x} \left[P(x) \frac{\partial u}{\partial x} \right] + Q(x)u - R(x), \quad u(-1) = \alpha, \quad u(1) = \beta.$$

In this case, the matrix equation we need to solve for the amplitudes \mathbf{a} of the basis functions is

$$C \frac{d\mathbf{a}}{dt} + K\mathbf{a} = \mathbf{f}.$$

This can be re-arranged to give:

$$\frac{d\mathbf{a}}{dt} = A\mathbf{a} + \mathbf{g}, \quad A = -C^{-1}K, \quad \mathbf{g} = C^{-1}\mathbf{f}.$$

We solve this matrix ODE by introducing a time lattice $t_i = ih$, where h is the stepsize. We write $\mathbf{a}_i = \mathbf{a}(t_i)$ and use a trapezoidal stencil to generate the numeric solution:

$$\left(I - \frac{h}{2}A\right)\mathbf{a}_{i+1} = \left(I + \frac{h}{2}A\right)\mathbf{a}_i + h\mathbf{g}$$

Here is a procedure that implements this method assuming initial data of the form $u(0, x) = \kappa(x)$. $\mathbf{N1}$ is the number of interior finite elements, $\mathbf{N2}$ is the number of time steps, and we assume the simulation runs from $t=0$ to $t=t_{\max}$:

```
> ParabolicSolver := proc(N1, N2, kappa, t_max, lambda, P, Q, R, alpha, beta)
  local phi, X, T, a, K, C, f, Cinv, A, g, h, Gamma, p, i, N:

  phi := GenerateBasis(N1):
  X := i -> -1 + 2*i/(N1+1):
  T := j -> j/N2*t_max:
  a := Vector([seq(evalf(kappa(X(i))), i=1..N1)], datatype=float):
  K := stiffness(phi, P, Q):
  C := damping(phi, lambda):
  f := force(phi, P, Q, R, alpha, beta):
  Cinv := C^(-1):
  A := -Cinv.K:
  g := Cinv.f:
  h := evalf(T(1)-T(0)):
  Gamma[1] := 1 - h*A/2:
  Gamma[2] := 1 + h*A/2:
  p[0] := Frame(a, phi, alpha, beta, T(0)):
  for i from 1 to N2 do
    a := LinearSolve(Gamma[1], Gamma[2].a+h*g):
```

```

    p[i] := Frame(a,phi,alpha,beta,T(i)):
od:
display(convert(p,list),insequence=true):

end proc:

Frame := proc(ans,Phi,alpha,beta,T) local N:
    N := rhs(ArrayDims(Phi))-1;
    plot([seq(alpha*Phi[0](x)+add(ans[i]*Phi[i](x),i=1..N)+beta*Phi[N+1](x),j=1..2)],x=-1.
.1,filled=[true,false],color=["Lavender","Purple"],axes=boxed,title=typeset(t=evalf[4](T)
),labels=[x,u(t,x)]);
end proc:

```

Here is an example of the output of **ParabolicSolver**:

```

> P := x -> 1-x^2:
Q := x -> 0:
R := x -> 0:
alpha := -1;
beta := 1;
rho := x -> 0:
lambda := x -> 1:
pde;
kappa := x -> sin(5*Pi*x/2);
N1 := 20:
N2 := 100:
t_max := 1.5:
ParabolicSolver(N1,N2,kappa,t_max,lambda,P,Q,R,alpha,beta);

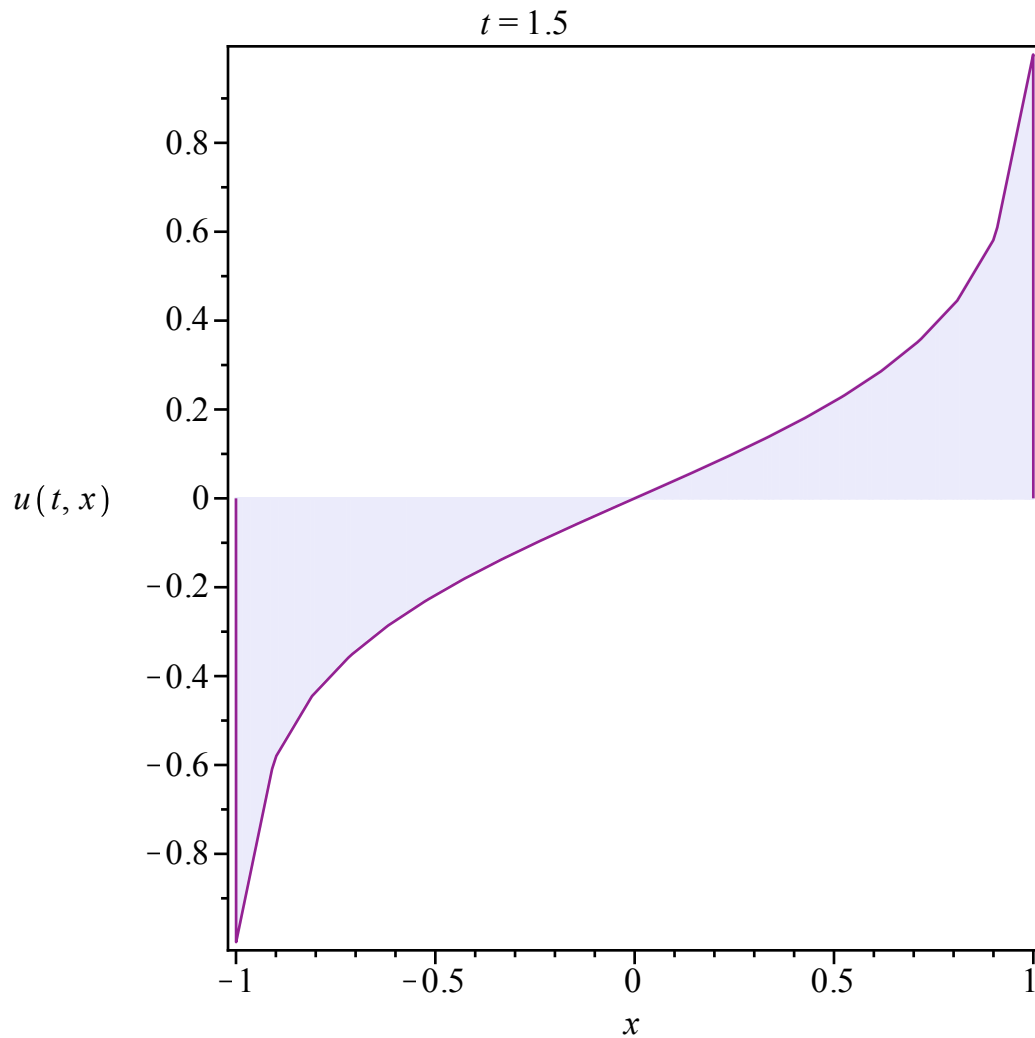
```

$$\alpha := -1$$

$$\beta := 1$$

$$\frac{\partial}{\partial t} u(t,x) - \left(\frac{\partial}{\partial x} \left((1-x^2) \left(\frac{\partial}{\partial x} u(t,x) \right) \right) \right)$$

$$\kappa := x \rightarrow \sin\left(\frac{5}{2} \pi x\right)$$



▼ Hyperbolic initial value problems

In this section, we concentrate on the complete PDE (1.1):

$$\rho(x) \frac{\partial^2 u}{\partial t^2} + \lambda(x) \frac{\partial u}{\partial t} = \frac{\partial}{\partial x} \left[P(x) \frac{\partial u}{\partial x} \right] + Q(x)u - R(x), \quad u(t, -1) = \alpha, \quad u(t, 1) = \beta.$$

In this case, the matrix ODE to solve is

$$M \frac{d^2 \mathbf{a}}{dt^2} + C \frac{d\mathbf{a}}{dt} + K\mathbf{a} = \mathbf{f}.$$

We convert this into a set of coupled first order ODEs via the definition $\mathbf{b} = \frac{d\mathbf{a}}{dt}$:

$$\frac{d\mathbf{a}}{dt} = \mathbf{b}, \quad \frac{d\mathbf{b}}{dt} = -M^{-1}K\mathbf{a} - M^{-1}C\mathbf{b} + M^{-1}\mathbf{f},$$

This can in turn be represented by a single matrix ODE:

$$\frac{d\mathbf{B}}{dt} = A\mathbf{B} + \mathbf{g}, \quad \mathbf{B} = \begin{bmatrix} \mathbf{a} \\ \mathbf{b} \end{bmatrix}, \quad A = \begin{bmatrix} 0 & I \\ -M^{-1}K & -M^{-1}C \end{bmatrix}, \quad \mathbf{g} = \begin{bmatrix} 0 \\ M^{-1}\mathbf{f} \end{bmatrix}.$$

We will solve this with the same type of trapezoidal stencil that we used in the last section; i.e., we write $\mathbf{B}_i = \mathbf{B}(t_i)$ with $t_{i+1} - t_i = h$ and use a trapezoidal stencil to generate the numeric solution:

$$\left(I - \frac{h}{2}A \right) \mathbf{B}_{i+1} = \left(I + \frac{h}{2}A \right) \mathbf{B}_i + h\mathbf{g}$$

Here is a procedure that implements this method assuming initial data of the form $u(0, x) = \kappa(x)$ and $\dot{u}(0, x) = \xi(x)$. **N1** is the number of interior finite elements, **N2** is the number of time steps, and we assume the simulation runs from **t=0** to **t=t_max**:

```
> HyperbolicSolver := proc(N1, N2, kappa, xi, t_max, rho, lambda, P, Q, R, alpha, beta)
  local phi, X, T, a, K, C, f, Minv, A, g, h, Gamma, p, i, N, b, B, M:

  phi := GenerateBasis(N1):
  X := i -> -1 + 2*i/(N1+1):
  T := j -> j/N2*t_max:
  a := Vector([seq(evalf(kappa(X(i))), i=1..N1)], datatype=float):
  b := Vector([seq(evalf(xi(X(i))), i=1..N1)], datatype=float):
  B := Vector([a, b]):
  K := stiffness(phi, P, Q):
```

```

C := damping(phi, lambda) :
M := mass(phi, rho) :
f := force(phi, P, Q, R, alpha, beta) :
Minv := M^(-1) :
A := Matrix( [[ZeroMatrix(N1, N1), IdentityMatrix(N1, N1)], [-Minv.K, -Minv.C]] ) :
g := Vector( [ZeroMatrix(N1, 1), Minv.f] ) :
h := evalf(T(1)-T(0)) :
Gamma[1] := 1 - h*A/2 :
Gamma[2] := 1 + h*A/2 :
p[0] := Frame(B[1..N1], phi, alpha, beta, T(0)) :
for i from 1 to N2 do :
    B := LinearSolve(Gamma[1], Gamma[2].B+h*g) :
    p[i] := Frame(B[1..N1], phi, alpha, beta, T(i)) :
od :
display(convert(p, list), insequence=true) :

end proc :

```

Here is some example output:

```

> P := x -> 1-x^2 :
Q := x -> 0 :
R := x -> 0 :
alpha := -1 :
beta := 1 :
rho := x -> 1 :
lambda := x -> 0 :
pde :
kappa := x -> sin(Pi*x/2) :
xi := x -> x^2 :
N1 := 40 :
N2 := 100 :
t_max := 20 :
HyperbolicSolver(N1, N2, kappa, xi, t_max, rho, lambda, P, Q, R, alpha, beta) ;

```

$$\alpha := -1$$

$$\beta := 1$$

$$\frac{\partial^2}{\partial t^2} u(t, x) - \left(\frac{\partial}{\partial x} \left((1 - x^2) \left(\frac{\partial}{\partial x} u(t, x) \right) \right) \right)$$

$$\kappa := x \rightarrow \sin\left(\frac{1}{2} \pi x\right)$$

$$\xi := x \rightarrow x^2$$

$t = 20.$

